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RELATIVE ENTROPY METHOD FOR MEASURE SOLUTIONS OF THE GROWTH-FRAGMENTATION EQUATION

TOMASZ DEBIEC, MARIE DOUMIC, PIOTR GWIAZDA, AND EMIL WIEDEMANN

ABSTRACT. The aim of this study is to generalise recent results of the two last authors on entropy methods for measure solutions of the renewal equation to other classes of structured population problems. Specifically, we develop a generalised relative entropy inequality for the growth-fragmentation equation and prove asymptotic convergence to a steady-state solution, even when the initial datum is only a non-negative measure.

1. INTRODUCTION

Structured population models were developed for the purpose of understanding the evolution of a population over time - and in particular to adequately describe the dynamics of a population by its distribution along some "structuring" variables representing e.g., age, size, or cell maturity. These models, often taking the form of an evolutionary partial differential equation, have been extensively studied for many years. The first age structure was considered in the early 20th century by Sharpe and Lotka [34], who already made predictions on the question of asymptotic behaviour of the population, see also [23, 24]. In the second half of the 20th century size-structured models appeared first in [3, 35]. These studies gave rise to other physiologically structured models (age-size, saturation, cell maturity, etc.).

The object of this note is the growth-fragmentation model, which is found fitting in many different contexts: cell division, polymerisation, neurosciences, prion proliferation or even telecommunication. In its general linear form this model takes the form of the following equation.

$$\begin{aligned}\partial_t n(t, x) + \partial_x(g(x)n(t, x)) + B(x)n(t, x) &= \int_x^\infty k(x, y)B(y)n(t, y) dy, \\ g(0)n(t, 0) &= 0, \\ n(0, x) &= n^0(x).\end{aligned}\tag{1.1}$$

Here $n(t, x)$ represents the concentration of individuals of size $x \geq 0$ at time $t > 0$, $g(x) \geq 0$ is their growth rate, $B(x) \geq 0$ is their division rate and $k(x, y)$ is the proportion of individuals of size x created out of the division of individuals of size y . This equation incorporates a very important phenomenon in biology - a competition between growth and fragmentation. Clearly they have opposite dynamics: growth drives the population towards a larger size, while fragmentation makes it smaller and smaller. Depending on which factor dominates, one can observe various long-time behaviour of the population distribution.

Many authors have studied the long-time asymptotics (along with well-posedness) of variants of the growth-fragmentation equation, see e.g. [7, 15, 27, 30, 33]. The studies which establish convergence, in a proper sense, of a (renormalised) solution towards a steady profile were until recently limited only to initial data in weighted L^1 spaces. The classical tools for such studies include a direct application of the Laplace transform and the semigroup theory [30]. These methods could also provide an exponential rate of convergence, linked to the existence of a spectral gap.

Key words and phrases. measure solutions, growth-fragmentation equation, structured population, relative entropy, generalised Young measure.

A different approach was developed by Perthame et al. in a series of papers [28, 29, 33]. Their Generalised Relative Entropy (GRE) method provides a way to study long-time asymptotics of linear models even when no spectral gap is guaranteed - however failing to provide a rate of convergence, unless an entropy-entropy dissipation inequality is obtained [7]. Recently Gwiazda and Wiedemann [21] extended the GRE method to the case of the renewal equation with initial data in the space of non-negative Radon measures. Their result is motivated by the increasing interest in measure solutions to models of mathematical biology, see e.g. [8, 20] for some recent results concerning well-posedness and stability theory in the space of non-negative Radon measures. The clear advantage of considering measure data is that it is biologically justified - it allows for treating the situation when a population is initially concentrated with respect to the structuring variable (and is, in particular, not absolutely continuous with respect to the Lebesgue measure). This is typically the case when departing from a population formed by a unique cell. We refer also to the recent result of Gabriel [18], who uses the Doeblin method to analyze the long-time behaviour of measure solutions to the renewal equation.

Let us remark that the method of analysis employed in the current paper is inspired by the classical relative entropy method introduced by Dafermos in [10]. In recent years this method was extended to yield results on measure-valued-strong uniqueness for equations of fluid dynamics [6, 17, 22] and general conservation laws [9, 12, 19]. See also [11] and references therein.

The purpose of this paper is to generalise the results of [21] to the case of a general growth-fragmentation equation. Similarly as in that paper we make use of the concept of a recession function to make sense of compositions of nonlinear functions with a Radon measure. However, the appearance of the term $H'(u_\varepsilon(t, x))u_\varepsilon(t, y)$ in the entropy dissipation (see (3.8) below), which mixes dependences on the variables x and y , poses a novel problem, which is overcome by using generalised Young measures and time regularity.

The current paper is structured as follows: in Section 2 we recall some basic results on Radon measures, recession functions and Young measures as well as introduce the assumptions of our model, in Section 3 we state and prove the GRE inequality, which is then used to prove a long-time asymptotics result in Section 4.

2. DESCRIPTION OF THE MODEL

2.1. Preliminaries. In what follows we denote by \mathbb{R}_+ the set $[0, \infty)$. By $\mathcal{M}(\mathbb{R}_+)$ we denote the space of signed Radon measures on \mathbb{R}_+ . By Lebesgue's decomposition theorem, for each $\mu \in \mathcal{M}(\mathbb{R}_+)$ we can write

$$\mu = \mu^a + \mu^s,$$

where μ^a is absolutely continuous with respect to the Lebesgue measure \mathcal{L}^1 , and μ^s is singular. The space $\mathcal{M}(\mathbb{R}_+)$ is endowed with the total variation norm

$$\|\mu\|_{\text{TV}} := \int_{\mathbb{R}_+} d|\mu|,$$

and we denote $\|\mu\|_{\text{TV}} = TV(\mu)$. By the Riesz Representation Theorem we can identify this space with the dual space to the space $\mathcal{C}_0(\mathbb{R}_+)$ of continuous functions on \mathbb{R}_+ which vanish at infinity. The duality pairing is given by

$$\langle \nu, f \rangle := \int_0^\infty f(\xi) d\mu(\xi).$$

By $\mathcal{M}^+(\mathbb{R}_+)$ we denote the set of positive Radon measures of bounded total variation. We further define the φ -weighted total variation by

$$\|\mu\|_{TV_\varphi} := \int_{\mathbb{R}_+} \varphi d|\mu|$$

and correspondingly the space $\mathcal{M}^+(\mathbb{R}_+; \varphi)$ of positive Radon measures whose φ -weighted total variation is finite. We still denote $TV(\mu) = \|\mu\|_{TV_\varphi}$. Of course we require that the function φ be non-negative. In fact, for our purposes φ will be strictly positive and bounded on each compact subset of $(0, \infty)$.

We say that a sequence $\nu_n \in \mathcal{M}(\mathbb{R}_+)$ converges *weak** to some measure $\nu \in \mathcal{M}(\mathbb{R}_+)$ if

$$\langle \nu_n, f \rangle \longrightarrow \langle \nu, f \rangle$$

for each $f \in \mathcal{C}_0(\mathbb{R}_+)$.

By a *Young measure* on $\mathbb{R}_+ \times \mathbb{R}_+$ we mean a parameterised family $\nu_{t,x}$ of probability measures on \mathbb{R}_+ . More precisely, it is a weak*-measurable function $(t, x) \mapsto \nu_{t,x}$, i.e. such that the mapping

$$(t, x) \mapsto \langle \nu_{t,x}, f \rangle$$

is measurable for each $f \in C_0(\mathbb{R}_+)$. Young measures are often used to describe limits of weakly converging approximating sequences to a given problem. They serve as a way of describing weak limits of nonlinear functions of the approximate solution. Indeed, it is a classical result that a uniformly bounded measurable sequence u_n generates a Young measure by which one represents the limit of $f(u_n)$, where f is some non-linear function, see [?] for sequences in L^∞ and [2] for measurable sequences.

This framework was used by DiPerna in his celebrated paper [13], where he introduced the concept of an admissible measure-valued solution to scalar conservation laws. However, in more general contexts (e.g. for hyperbolic systems, where there is usually only one entropy-entropy-flux pair) one needs to be able to describe limits of sequences which exhibit oscillatory behaviour as well as concentrate mass. Such a framework is provided by *generalised Young measures*, first introduced in the context of incompressible Euler equations in [14], and later developed by many authors. We follow the exposition of Alibert, Bouchitté [1] and Kristensen, Rindler [26].

Suppose $f : \mathbb{R}^n \rightarrow \mathbb{R}_+$ is an even continuous function with at most linear growth, i.e.

$$|f(x)| \leq C(1 + |x|)$$

for some constant C . We define, whenever it exists, the *recession function* of f as

$$f^\infty(x) = \lim_{s \rightarrow \infty} \frac{f(sx)}{s} = \lim_{s \rightarrow \infty} \frac{f(-sx)}{s}.$$

Definition 2.1. The set $\mathcal{F}(\mathbb{R})$ of continuous functions $f : \mathbb{R} \rightarrow \mathbb{R}_+$ for which f^∞ exists and is continuous on \mathbb{S}^{n-1} is called the class of admissible integrands.

By a *generalised Young measure* on $\Omega = \mathbb{R}_+ \times \mathbb{R}_+$ we mean a parameterised family $(\nu_{t,x}, m)$ where for $(t, x) \in \Omega$, $\nu_{t,x}$ is a family of probability measures on \mathbb{R} and m is a nonnegative Radon measure on Ω . In the following, we may omit the indices for $\nu_{t,x}$ and denote it simply (ν, m) .

The following result gives a way of representing weak* limits of sequences bounded in L^1 via a generalised Young measure. It was first proved in [1, Theorem 2.5]. We state an adaptation to our simpler case.

Proposition 2.2. Let (u_n) be a bounded sequence in $L^1_{loc}(\Omega; \mu, \mathbb{R})$, where μ is a measure on Ω . There exists a subsequence (u_{n_k}) , a nonnegative Radon measure m on Ω and a parametrized family

of probabilities (ν_ζ) such that for any even function $f \in \mathcal{F}(\mathbb{R})$ we have

$$f(u_{n_k}(\zeta))\mu \xrightarrow{*} \langle \nu_\zeta, f \rangle \mu + f^\infty m \quad (2.1)$$

Proof. We apply Theorem 2.5. and Remark 2.6 in [1], simplified by the fact that f is even and that we only test against functions f independent of x . Note that the weak* convergence then has to be understood in the sense of compactly supported test functions $\varphi \in \mathcal{C}_0(\Omega, \mathbb{R})$. \square

The above proposition can in fact be generalised to say that every bounded sequence of generalised Young measures possesses a weak* convergent subsequence, cf. [26, Corollary 2.]

Proposition 2.3. *Let (ν^n, m^n) be a sequence of generalised Young measures on Ω such that*

- *The map $x \mapsto \langle \nu_x^n, |\cdot| \rangle$ is uniformly bounded in L^1 ,*
- *The sequence $(m^n(\bar{\Omega}))$ is uniformly bounded.*

Then there is a generalised Young measure (ν, m) on Ω such that (ν^n, m^n) converges weak to (ν, m) .*

2.2. The model. We consider the growth-fragmentation equation under a general form:

$$\begin{aligned} \partial_t n(t, x) + \partial_x(g(x)n(t, x)) + B(x)n(t, x) &= \int_x^\infty k(x, y)B(y)n(t, y) dy, \\ g(0)n(t, 0) &= 0, \\ n(0, x) &= n^0(x). \end{aligned} \quad (2.2)$$

We assume $n^0 \in \mathcal{M}^+(\mathbb{R}_+)$.

The fundamental tool in studying the long-time asymptotics with the GRE method is the existence and uniqueness of the first eigenelements (λ, N, φ) , i.e. solutions to the following primal and dual eigenproblems.

$$\begin{aligned} \frac{\partial}{\partial x}(g(x)N(x)) + (B(x) + \lambda)N(x) &= \int_x^\infty k(x, y)B(y)N(y) dy \\ g(0)N(0) &= 0, \quad N(x) > 0, \text{ for } x > 0, \quad \int_0^\infty N(x)dx = 1, \end{aligned} \quad (2.3)$$

$$\begin{aligned} -g(x)\frac{\partial}{\partial x}(\varphi(x)) + (B(x) + \lambda)\varphi(x) &= B(x)\int_0^x k(y, x)\varphi(y) dy \\ \varphi(x) &> 0, \quad \int_0^\infty \varphi(x)N(x)dx = 1. \end{aligned} \quad (2.4)$$

We make the following assumptions on the parameters of the model.

$$B \in W^{1,\infty}(\mathbb{R}_+, \mathbb{R}_+^*), \quad g \in W^{1,\infty}(\mathbb{R}_+, \mathbb{R}_+^*), \quad \forall x \geq 0, \quad g \geq g_0 > 0, \quad (2.5)$$

$$k \in \mathcal{C}_b(\mathbb{R}_+ \times \mathbb{R}_+), \quad \int_0^y k(x, y)dx = 2, \quad \int_0^y xk(x, y)dx = y, \quad (2.6)$$

$$k(x, y < x) = 0, \quad k(x, y > x) > 0. \quad (2.7)$$

These guarantee in particular existence and uniqueness of a solution $n \in \mathcal{C}(\mathbb{R}_+; L_\varphi^1(\mathbb{R}_+))$ for L^1 initial data (see e.g. [31]), existence of a unique measure solution for data in $\mathcal{M}^+(\mathbb{R}_+)$ (cf. [8]), as well as existence and uniqueness of a dominant eigentriplet $(\lambda > 0, N(x), \varphi(x))$, cf. [15]. In particular the functions N and φ are continuous, N is bounded and φ has at most polynomial growth. In what follows N and φ will always denote the solutions to problems (2.3) and (2.4), respectively. Let us remark that in the L^1 setting we have the following conservation law

$$\int_0^\infty n_\varepsilon(t, x)e^{-\lambda t}\varphi(x)dx = \int_0^\infty n^0(x)\varphi(x)dx. \quad (2.8)$$

2.3. Measure and measure-valued solutions. Let us observe that there are two basic ways to treat the above model in the measure setting. The first one is to consider a *measure solution*, i.e. a narrowly continuous map $t \mapsto \mu_t \in \mathcal{M}^+(\mathbb{R}_+)$, which satisfies (2.2) in the weak sense, i.e. for each $\psi \in \mathcal{C}_c^1(\mathbb{R}_+ \times \mathbb{R}_+)$

$$\begin{aligned} & - \int_0^\infty \int_0^\infty (\partial_t \psi(t, x) + \partial_x \psi(t, x) g(x)) d\mu_t(x) dt + \int_0^\infty \int_0^\infty \psi(t, x) B(x) d\mu_t(x) dt \\ & = \int_0^\infty \int_0^\infty \psi(t, x) \int_x^\infty k(x, y) B(y) d\mu_t(y) dx dt + \int_0^\infty \psi(0, x) dn^0(x). \end{aligned} \quad (2.9)$$

Thus a measure solution is a family of time-parameterised non-negative Radon measures on the structure-physical domain \mathbb{R}_+ .

The second way is to work with generalised Young measures and corresponding measure-valued solutions. To prove the generalised relative entropy inequality, which relies on considering a family of non-linear renormalisations of the equation, we choose to work in this second framework.

A *measure-valued solution* is a generalised Young measure (ν, m) , where the oscillation measure is a family of parameterised probabilities over the state domain \mathbb{R}_+ such that equation (2.2) is satisfied by its barycenters $\langle \nu_{t,x}, \xi \rangle$, i.e. the following equation

$$\begin{aligned} & \partial_t (\langle \nu_{t,x}, \xi \rangle + m) + \partial_x (g(x) (\langle \nu_{t,x}, \xi \rangle + m)) + B(x) (\langle \nu_{t,x}, \xi \rangle + m) \\ & = \int_x^\infty k(x, y) B(y) \langle \nu_{t,x}, \xi \rangle dy + \int_x^\infty k(x, y) B(y) dm(y) \end{aligned} \quad (2.10)$$

holds in the sense of distributions on $\mathbb{R}_+^* \times \mathbb{R}_+^*$.

It is proven in [20] that equation (2.2) has a unique measure solution. To this solution one can associate a measure-valued solution - for example, given a measure solution $t \mapsto \mu_t$ one can define a measure-valued solution by

$$\langle \delta_{\left\{ \frac{d\mu_t^a}{d\mu_1} \right\}}, \text{id} \rangle = \mu_t^a, \quad m = \mu_t^s$$

where $\frac{d\mu_1}{d\mu_2}$ denotes the Radon-Nikodym derivative of μ_1 with respect to μ_2 .

However, clearly, the measure-valued solutions are not unique - since the equation is linear, there is freedom in choosing the Young measure as long as the barycenter satisfies equation (2.10). For example, a different measure-valued solution can be defined by

$$\left\langle \frac{1}{2} \delta_{\left\{ 2 \frac{d\mu_t^a}{d\mu_1} \right\}} + \frac{1}{2} \delta_{\{0\}}, \text{id} \right\rangle = \mu_t^a.$$

Uniqueness of measure-valued solution can be ensured by requiring that the generalised Young measure satisfies not only the equation, but also a family of nonlinear renormalisations. This was the case in the work of DiPerna [13], see also [11].

To establish the GRE inequality which will then be used to prove an asymptotic convergence result, we consider the measure-valued solution generated by a sequence of regularized solutions. This allows us to use the classical GRE method established in [32]. Careful passage to the limit will then show that analogous inequalities hold for the measure-valued solution.

3. GRE INEQUALITY

In this section we formulate and prove the generalised relative entropy inequality, our main tool in the study of long-time asymptotics for equation (2.2). We take advantage of the well-known GRE inequalities in the L^1 setting. To do so we consider the growth-fragmentation equation (2.2) for a sequence of regularized data and prove that we can pass to the limit, thus obtaining the desired inequalities in the measure setting.

Let $n_\varepsilon^0 \in L_\varphi^1(\mathbb{R}_+)$ be a sequence of regularizations of n^0 converging weak* to n^0 in the space of measures and such that $\text{TV}(n_\varepsilon^0) \rightarrow \text{TV}(n^0)$. Let n_ε denote the corresponding unique solution to (2.2) with n_ε^0 as an initial condition. Then for any differentiable strictly convex admissible integrand H we define the usual relative entropy

$$\mathcal{H}_\varepsilon(t) := \int_0^\infty \varphi(x) N(x) H\left(\frac{n_\varepsilon(t, x) e^{-\lambda t}}{N(x)}\right) dx$$

and entropy dissipation

$$\begin{aligned} D_\varepsilon^H(t) = \int_0^\infty \int_0^\infty \varphi(x) N(y) B(y) k(x, y) & \left\{ H\left(\frac{n_\varepsilon(t, y) e^{-\lambda t}}{N(y)}\right) - H\left(\frac{n_\varepsilon(t, x) e^{-\lambda t}}{N(x)}\right) \right. \\ & \left. - H'\left(\frac{n_\varepsilon(t, x) e^{-\lambda t}}{N(x)}\right) \left[\frac{n_\varepsilon(t, y) e^{-\lambda t}}{N(y)} - \frac{n_\varepsilon(t, x) e^{-\lambda t}}{N(x)} \right] \right\} dx dy. \end{aligned}$$

Then, as shown e.g. in [28], one can show that

$$\frac{d}{dt} \left\{ \int_0^\infty \varphi(x) N(x) H\left(\frac{n_\varepsilon(t, x) e^{-\lambda t}}{N(x)}\right) dx \right\} = -D_\varepsilon^H(t) \quad (3.1)$$

with the right-hand side being non-positive due to convexity of H . Hence the relative entropy is non-increasing. It follows that $\mathcal{H}_\varepsilon(t) \leq \mathcal{H}_\varepsilon(0)$ and, since $H \geq 0$,

$$\int_0^\infty D_\varepsilon^H(t) dt \leq \mathcal{H}_\varepsilon(0). \quad (3.2)$$

In the next proposition we prove corresponding inequalities for the measure-valued solution generated by the sequence n_ε . This result is an analogue of Theorem 5.1 in [21].

Proposition 3.1. *With notation as above, there exists a subsequence (not relabelled), generating a generalised Young measure (ν, m) with $m = m_t \otimes dt$ for a family of positive Radon measures m_t , such that*

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_0^\infty \chi(t) \mathcal{H}_\varepsilon(t) dt &= \int_0^\infty \chi(t) \left(\int_0^\infty \varphi(x) N(x) \langle \nu_{t,x}(\alpha), H(\alpha) \rangle dx \right. \\ &\quad \left. + \int_0^\infty \varphi(x) N(x) H^\infty dm_t(x) \right) dt \end{aligned} \quad (3.3)$$

for any $\chi \in C_c([0, \infty))$, and

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_0^\infty D_\varepsilon^H(t) dt &= \\ \int_0^\infty \int_0^\infty \int_0^\infty \varphi(x) N(y) B(y) k(x, y) & \langle \nu_{t,y}(\xi) \otimes \nu_{t,x}(\alpha), H(\xi) - H(\alpha) - H'(\alpha)(\xi - \alpha) \rangle dx dy dt \\ + \int_0^\infty \int_0^\infty \int_0^\infty \varphi(x) N(y) B(y) k(x, y) & \langle \nu_{t,x}(\alpha), H^\infty - H'(\alpha) \rangle dm_t(y) dx dt \geq 0. \end{aligned} \quad (3.4)$$

We denote the limits on the left-hand sides of the above equations by $\int_0^\infty \chi(t) \mathcal{H}(t) dt$ and $\int_0^\infty D^H(t) dt$, respectively, thus defining the measure-valued relative entropy and entropy dissipation for almost every t . We further set

$$\mathcal{H}(0) := \int_0^\infty \varphi(x) N(x) H\left(\frac{(n^0)^a(x)}{N(x)}\right) dx + \int_0^\infty \varphi(x) H^\infty\left(\frac{(n^0)^s(x)}{|(n^0)^s|}\right) d|(n^0)^s(x)|. \quad (3.5)$$

We then have

$$\frac{d}{dt} \mathcal{H}(t) \leq 0 \quad \text{in the sense of distributions,} \quad (3.6)$$

and

$$\int_0^\infty D^H(t) dt \leq \mathcal{H}(0). \quad (3.7)$$

Proof. The function $t \mapsto \int_0^\infty n_\varepsilon(t, x) e^{-\lambda t} \varphi(x) dx$ is constant and the function N is strictly positive on $(0, \infty)$. Therefore the sequence $u_\varepsilon(t, x) := \frac{n_\varepsilon(t, x) e^{-\lambda t}}{N(x)}$ is uniformly bounded in $L^\infty(\mathbb{R}_+; L^1_{\varphi, \text{loc}}(\mathbb{R}_+))$. Hence we can apply Proposition 2.2 to obtain a generalised Young measure (ν, m) on $\mathbb{R}_+ \times \mathbb{R}_+$.

Since $u_\varepsilon \in L^\infty(\mathbb{R}_+; L^1_{\varphi, \text{loc}}(\mathbb{R}_+))$, we have $m \in L^\infty(\mathbb{R}_+; \mathcal{M}(\mathbb{R}_+; \varphi))$. By a standard disintegration argument (see for instance [16, Theorem 1.5.1]) we can write the slicing measure for m , $m(dt, dx) = m_t(dx) \otimes dt$, where the map $t \mapsto m_t \in \mathcal{M}^+(\mathbb{R}_+; \varphi)$ is measurable and bounded.

By Proposition 2.2 we have the weak* convergence

$$H(u_\varepsilon(t, x))(dt \otimes \varphi(x) dx) \xrightarrow{*} \langle \nu_{t, x}, H \rangle (dt \otimes \varphi(x) dx) + H^\infty m.$$

Testing with $(t, x) \mapsto \chi(t) N(x)$ where $\chi \in \mathcal{C}_c(\mathbb{R}_+)$, we obtain (3.3). Further, the convergence $\int_0^\infty \chi(t) \mathcal{H}_\varepsilon(t) dt \rightarrow \int_0^\infty \chi(t) \mathcal{H}(t) dt$ implies (3.6), since for \mathcal{H}_ε we have the corresponding inequality (3.1).

We now investigate the limit as $\varepsilon \rightarrow 0$ of $\int_0^\infty D_\varepsilon^H(t) dt$. Denoting $\Phi(x, y) := k(x, y) N(y) B(y)$ we have

$$D_\varepsilon^H(t) = \int_0^\infty \int_0^\infty \Phi(x, y) \varphi(x) [H(u_\varepsilon(t, y)) - H(u_\varepsilon(t, x)) - H'(u_\varepsilon(t, x)) u_\varepsilon(t, y) + H'(u_\varepsilon(t, x)) u_\varepsilon(t, x)] dx dy. \quad (3.8)$$

We consider each of the four terms of the sum separately on the restricted domain $[0, T] \times [\eta, K]^2$ for fixed $T > 0$ and $K > \eta > 0$. Let $D_{\varepsilon, \eta, K}^H$ denote the entropy dissipation with the integrals of (3.8) each taken over the subsets $[\eta, K]$ of \mathbb{R}_+ .

We now apply Proposition 2.2 to the sequence u_ε , the measure $dt \otimes \varphi(x) dx$ on the set $[0, T] \times [\eta, K]$. The first two and the last integrands of $D_{\varepsilon, \eta, K}^H(t)$ depend on t and only either on x or on y . Therefore we can pass to the limit as $\varepsilon \rightarrow 0$ by Proposition 2.2 using a convenient test function. More precisely, testing with $(t, x) \mapsto \int_\eta^K \Phi(x, y) dy$, we obtain the convergence

$$\begin{aligned} - \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) H(u_\varepsilon(t, x)) dy dx dt &\longrightarrow - \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) \langle \nu_{t, x}, H \rangle dy dx dt \\ &\quad - \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) H^\infty dm_t(x) dy dx dt, \end{aligned}$$

and, noticing that the recession function of $\alpha \mapsto \alpha H'(\alpha)$ is H^∞ ,

$$\begin{aligned} \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) H'(u_\varepsilon(t, x)) u_\varepsilon(t, x) dy dx dt &\longrightarrow \\ \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) \langle \nu_{t, x}, \alpha H'(\alpha) \rangle dy dx dt &+ \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) H^\infty dm_t(x) dy dx dt. \end{aligned}$$

Likewise, using $(t, y) \mapsto \frac{1}{\varphi(y)} \int_\eta^K \Phi(x, y) \varphi(x) dx$, we obtain

$$\begin{aligned} \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) H(u_\varepsilon(t, y)) dx dy dt &\longrightarrow \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) \langle \nu_{t, y}, H \rangle dx dy dt \\ &\quad + \int_0^T \int_\eta^K \int_\eta^K \Phi(x, y) \varphi(x) H^\infty dm_t(y) dx dy dt. \end{aligned}$$

There remains the term of $D_{\varepsilon,\eta,K}^H$ in which the dependence on u_ε combines x and y . To deal with this term we separate variables by testing against functions of the form $f_1(x)f_2(y)$. We then consider

$$\begin{aligned} & - \int_0^T \iint_{[\eta,K]^2} f_1(x)f_2(y)H'(u_\varepsilon(t,x))u_\varepsilon(t,y)dx dy dt \\ & = - \int_0^T \left(\int_\eta^K f_1(x)H'(u_\varepsilon(t,x))dx \right) \left(\int_\eta^K f_2(y)u_\varepsilon(t,y)dy \right) dt. \end{aligned}$$

The integrands are now split, one containing the x dependence, and one the y dependence. However, extra care is required here to pass to the limit. As the Young measures depend both on time and space, it is possible for the oscillations to appear in both directions. We therefore require appropriate time regularity of at least one of the sequences to guarantee the desired behaviour of the limit of the product.

Such requirement is met by noticing that since $u_\varepsilon \in \mathcal{C}([0,T];L_\phi^1(\mathbb{R}_+))$ uniformly, we have u_ε uniformly in $W^{1,\infty}([0,T];(\mathcal{M}^+(\mathbb{R}_+), \|\cdot\|_{(W^{1,\infty})^*}))$, cf. [8, Lemma 4.1]. Assuming $f_2 \in W^{1,\infty}(\mathbb{R}_+)$ we therefore have

$$\left(t \mapsto \int_\eta^K f_2(y)u_\varepsilon(t,y)dy \right) \in W^{1,\infty}([0,T]).$$

This in turn implies strong convergence of $\int_\eta^K f_2(y)u_\varepsilon(t,y)dy$ in $\mathcal{C}([0,T])$, by virtue of Arzela-Ascoli theorem. Therefore we have (noting that $(H')^\infty \equiv 0$ by sublinear growth of H)

$$\begin{aligned} & - \int_0^T \iint_{[\eta,K]^2} f_1(x)f_2(y)H'(u_\varepsilon(t,x))u_\varepsilon(t,y)dx dy dt \\ & = - \int_0^T \left(\int_\eta^K f_1(x)H'(u_\varepsilon(t,x))dx \right) \left(\int_\eta^K f_2(y)u_\varepsilon(t,y)dy \right) dt \\ & \longrightarrow - \int_0^T \left(\int_\eta^K f_1(x)\langle v_{t,x}, H' \rangle dx \right) \left(\int_\eta^K f_2(y)\langle v_{t,y}, \text{id} \rangle dy \right) dt \\ & \quad - \int_0^T \left(\int_\eta^K f_1(x)\langle v_{t,x}, H' \rangle dx \right) \left(\int_\eta^K f_2(y)dm_t(y) \right) dt \\ & = - \int_0^T \iint_{[\eta,K]^2} f_1(x)f_2(y)\langle v_{t,x}, H'(\alpha) \rangle \langle v_{t,y}, \xi \rangle dx dy \\ & \quad - \int_0^T \iint_{[\eta,K]^2} f_1(x)f_2(y)\langle v_{t,x}, H'(\alpha) \rangle dm_t(y) dx dt. \end{aligned}$$

By density of the linear space spanned by separable functions in the space of bounded continuous functions of (x,y) we obtain

$$\begin{aligned} & - \int_0^T \iint_{[\eta,K]^2} \Phi(x,y)\varphi(x)H'(u_\varepsilon(t,x))u_\varepsilon(t,y)dx dy dt \\ & \longrightarrow \int_0^T \iint_{[\eta,K]^2} \Phi(x,y)\varphi(x)\langle v_{t,x}, H'(\alpha) \rangle \langle v_{t,y}, \xi \rangle dx dy dt \\ & \quad - \int_0^T \iint_{[\eta,K]^2} \Phi(x,y)\varphi(x)\langle v_{t,x}, H'(\alpha) \rangle dm_t(y) dx dt. \end{aligned}$$

Gathering all the terms we thus obtain the convergence as $\varepsilon \rightarrow 0$

$$\int_0^T D_{\varepsilon,\eta,K}^H(t)dt \longrightarrow \int_0^T D_{\eta,K}^H(t)dt$$

with

$$\begin{aligned} D_{\eta,K}^H(t) &:= \iint_{[\eta,K]^2} \Phi(x,y) \varphi(x) \langle v_{t,y}(\xi) \otimes v_{t,x}(\alpha), H(\xi) - H(\alpha) - H'(\alpha)(\xi - \alpha) \rangle dx dy \\ &\quad + \iint_{[\eta,K]^2} \Phi(x,y) \varphi(x) \langle v_{t,x}(\alpha), H^\infty - H'(\alpha) \rangle dm_t(y) dx. \end{aligned}$$

Observe that since Φ is non-negative and H is convex, the integrand of $D_{\varepsilon,\eta,K}^H$ is non-negative. Hence so is the integrand of the limit. Therefore, by Monotone Convergence, we can pass to the limit $\eta \rightarrow 0$, $K \rightarrow \infty$, and $T \rightarrow \infty$ to obtain

$$\begin{aligned} 0 &\leq \lim_{\varepsilon \rightarrow 0} \int_0^\infty D_\varepsilon^H(t) dt = \\ &\int_0^\infty \int_0^\infty \int_0^\infty \varphi(x) N(y) B(y) k(x,y) \langle v_{t,y}(\xi) \otimes v_{t,x}(\alpha), H(\xi) - H(\alpha) - H'(\alpha)(\xi - \alpha) \rangle dx dy dt \\ &\quad + \int_0^\infty \int_0^\infty \int_0^\infty \varphi(x) N(y) B(y) k(x,y) \langle v_{t,x}(\alpha), H^\infty - H'(\alpha) \rangle dm_t(y) dx dt. \end{aligned}$$

Finally we note that by the Reshetnyak continuity theorem, cf. [21, 25] we have the convergence $\mathcal{H}_\varepsilon(0) \rightarrow \mathcal{H}(0)$. Together with (3.2) this implies (3.6). \square

4. LONG-TIME ASYMPTOTICS

In this section we use the result of the previous section to prove that a measure-valued solution of (2.2) converges to the steady-state solution. More precisely we prove

Theorem 4.1. *Let $n^0 \in \mathcal{M}(\mathbb{R}_+)$ and let n solve the growth-fragmentation equation (2.2). Then*

$$\lim_{t \rightarrow \infty} \int_0^\infty \varphi(x) d|n(t,x) - m_0 N(x) \mathcal{L}^1| = 0 \quad (4.1)$$

where $m_0 := \int_0^\infty \varphi(x) dn^0(x)$ and \mathcal{L}^1 denotes the 1-dimensional Lebesgue measure.

Proof. From inequality (3.7) we see that D^H belongs to $L^1(\mathbb{R}_+)$. Therefore there exists a sequence of times $t_n \rightarrow \infty$ such that

$$\lim_{n \rightarrow \infty} D^H(t_n) = 0.$$

Consider the corresponding sequence of generalised Young measures $(v_{t_n,x}, m_{t_n})$. Thanks to the inequality $\mathcal{H}(t) \leq \mathcal{H}(0)$ this sequence is uniformly bounded in the sense that

$$\sup_n \left\{ \int_0^\infty \varphi(x) N(x) \langle v_{t_n,x}(\alpha), |\alpha| \rangle dx + \int_0^\infty \varphi(x) N(x) dm_{t_n}(x) \right\} < \infty. \quad (4.2)$$

Therefore by the compactness property of Proposition 2.3 there is a subsequence, not relabelled, and a generalised Young measure (\bar{v}_x, \bar{m}) such that

$$(v_{t_n,x}, m_{t_n}) \xrightarrow{*} (\bar{v}_x, \bar{m})$$

in the sense of measures. We now show that the corresponding "entropy dissipation"

$$\begin{aligned} D_\infty^H &:= \int_0^\infty \int_0^\infty \Phi(x,y) \varphi(x) \langle \bar{v}_y(\xi) \otimes \bar{v}_x(\alpha), H(\xi) - H(\alpha) - H'(\alpha)(\xi - \alpha) \rangle dx dy \\ &\quad + \int_0^\infty \int_0^\infty \Phi(x,y) \varphi(x) \langle \bar{v}_x(\alpha), H^\infty - H'(\alpha) \rangle d\bar{m}(y) dx \end{aligned} \quad (4.3)$$

is zero. To this end we argue that

$$D_\infty^H = \lim_{n \rightarrow \infty} D^H(t_n).$$

Indeed this follows by the same arguments as in the proof of Proposition 3.1. In fact now the "mixed" term poses no additional difficulty as there is no time integral. It therefore follows that

$$D_\infty^H = 0. \quad (4.4)$$

As H is convex, both integrands in (4.3) are non-negative. Therefore (4.4) implies that both the integrals of D_∞^H are zero. In particular

$$\int_0^\infty \int_0^\infty H(\xi) - H(\alpha) - H'(\alpha)(\xi - \alpha) d\bar{\nu}_x(\alpha) d\bar{\nu}_y(\xi) = 0,$$

and since the integrand vanishes if and only if $\xi = \alpha$, this implies that the Young measure $\bar{\nu}$ is a Dirac measure concentrated at a constant. Then the vanishing of the second integral of D_∞^H implies that $\bar{m} = 0$. Moreover, the constant can be identified as

$$m_0 := \int_0^\infty \varphi(x) dn^0(x) \quad (4.5)$$

by virtue of the conservation in time of

$$\int_0^\infty \varphi(x) e^{-\lambda t} \langle v_{t,x}, \cdot \rangle dx + \int_0^\infty \varphi(x) e^{-\lambda t} dm_t(x).$$

By virtue of Proposition 2.2 with $H = |\cdot - m_0|$ it then follows that

$$\lim_{n \rightarrow \infty} \int_0^\infty \varphi(x) d|n(t_n, x) e^{-\lambda t_n} - m_0 N(x) \mathcal{L}^1| = 0,$$

which is the desired result, at least for our particular sequence of times.

Finally, we can argue that the last convergence holds for the entire time limit $t \rightarrow \infty$, invoking the monotonicity of the relative entropy \mathcal{H} . Indeed, the choice $H = |\cdot - m_0|$ in (3.5) yields the monotonicity in time of

$$\int_0^\infty \varphi(x) d|n(t, x) e^{-\lambda t} - m_0 N(x) \mathcal{L}^1|,$$

and the result follows. \square

CONCLUSION

In this article, we have proved the long-time convergence of measure-valued solutions to the growth-fragmentation equation. This result extends previously obtained results for L_φ^1 solutions [29]. As for the renewal equation [21], it is based on extending the generalised relative entropy inequality to measure-valued solutions, thanks to recession functions. Generalised Young measures provide an adequate framework to represent the measure-valued solutions and their entropy functionals.

Under slightly stronger assumptions on the fragmentation kernel k , e.g. the ones assumed in [7], it has been proved that an entropy-entropy dissipation inequality could be obtained. Under such assumptions, we could obtain in a simple way a stronger result of exponential convergence, see the proof of Theorem 4.1. in [21]. However the above seen method allows us to extend the convergence to spaces where no spectral gap exists [5].

A specially interesting case of application of this method would be critical cases where the dominant eigenvalue is not unique but is given by a countable set of eigenvalues. It has been proved that for L^2 initial conditions, the solution then converges to its projection on the space spanned by the dominant eigensolutions [4]. In the case of measure-valued initial condition, due to the fact that the equation has not anymore a regularisation effect, the asymptotic limit is expected to be the periodically oscillating measure, projection of the initial condition on the space of measures spanned by the dominant eigensolutions. This is a subject for future work.

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